

# Bioluminescence of Firefly Squid via Mechanism of Single Electron-Transfer Oxygenation and Charge-Transfer-Induced Luminescence

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# **Supporting Information**

**ABSTRACT:** Watasenia scintillans (W. scintillans) is a deepsea luminescent squid with a popular name of firefly squid. It produces flashes of blue light via a series of complicated luciferin-luciferase reactions involving ATP,  $Mg^{2+}$ , and molecular oxygen. Tsuji has proposed a hypothetical scheme for this mysterious bioluminescence (BL) process, but the proposal is short of strong evidence experimentally or theoretically, especially for two key steps. They are the addition of molecular oxygen to luciferin and the formation of light emitter. For the first time, the present study investigates the two steps by reliable density functional theory (DFT) and time-dependent DFT. The results of calculated energetics, charge transfer process, electronic structures, and molecular dynamics give convincing support for Tsuji's proposal. The



oxygenation reaction occurs with a single electron-transfer (SET) mechanism, and the light emitter is produced via the mechanism of gradually reversible charge-transfer-induced luminescence (GRCTIL). The simulation of nonadiabatic molecular dynamics further confirms the GRCTIL mechanisms and evaluates the quantum yield of the light emitter to be 43%. The knowledge obtained in the current study will help to understand a large amount of BL systems in nature, since the core structure of *W. scintillans* luciferin, imidazopyrazinone, is common in the luciferins of about eight phyla of luminescent organisms.

# 1. INTRODUCTION

Bioluminescence (BL) is an amazing phenomenon produced by a living organism. It has received extensive attention as a subject of life science and a powerful analytical tool in medical applications, such as noninvasive real-time imaging for target molecules in vivo.<sup>1-3</sup> There are more than 700 genera known to produce BL, and within them the most famous one could be the firefly.<sup>4</sup> Some other terrestrial creatures are also able to produce BL, such as fungi<sup>5</sup> and earthworms.<sup>6</sup> Nevertheless, about 80% of luminous organisms live in the ocean from bacteria to fish.<sup>7,8</sup> Watasenia scintillans (W. scintillans) is a luminescent squid living in the deep sea.<sup>9</sup> It has  $\sim 6$  cm mantle length and is almost entirely covered by thousands of photophores that emit light in synchronous or nonsynchronous pattern.<sup>10</sup> W. scintillans produces bright blue flashes (470 nm) like the firefly from its tips of arms, <sup>11</sup> so it is popularly called firefly squid. Firefly squids glow brilliantly off the coast when they gather in millions during spawning period.<sup>12</sup> How is the glamorous light produced? Generally speaking, luminous squid emits light in three ways.9 One, the light is emitted by symbiotic luminous bacteria harboring in its photophores, such as Hawaiian bobtail squid (Euprymna scolopes) that produces light by symbiotic bacteria (Vibrio fischeri).<sup>13–16</sup> Two, luminous mucus instead of ink is ejected from the luminous organ of squid into the sea. This kind of BL is common in deep-sea squid, such as Heteroteuthis dispar.<sup>17</sup> Three, the most common way for BL creatures is when the light is emitted by endocellular chemiluminescent reactions in photoprotein or luciferin-luciferase system. In the former case, a photoprotein is first formed, then the BL is produced after the photoprotein encounters some necessary chemicals. For instance, Symplectoteuthis oualaniensis emits light in the presence of oxygen molecule and metal cations (such as Na<sup>+</sup> and K<sup>+</sup>) after dehydrocoelenterazine (Chart 1A) links to a protein at a cysteine residue to form the photoprotein (symplectin, Chart 1B).<sup>18-23</sup> In the latter case, the luciferase catalyzes the oxidation of luciferin to produce the light emitter. The BL of W. scintillans is a typical example of the latter case.<sup>10,11</sup> The light is produced via an ATP-dependent luciferin-luciferase reaction (the optimal pH value of 8.8) involving oxygen molecule and  $Mg^{2+,9}$  The luciferin of W. scintillans is a disulfonate derivative of coelenterazine (Chart 1C), i.e., coelenterazine disulfonate (Chart 1D), which has been isolated from the photophores of W. scintillans.<sup>24-26</sup> Although coelenterazine was isolated in the liver of W. scintillans, coelenterazine disulfonate was not isolated in liver but in arm photophores under complete exclusion of oxygen, which gave

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<sup>*a*</sup>A is dehydrocoelenterazine, B is the photoprotein for *S. oualaniensis* BL (symplectin), C is coelenterazine, and D is the luciferin of *W. scintillans* (coelenterazine disulfonate), where the skeleton in blue (ImPy) is the common core of many luminescent substrates.





Scheme 2. <sup>3</sup>O<sub>2</sub> Addition Paths for ImPys at Three (C<sub>2</sub>-, C<sub>3</sub>-, or C<sub>5</sub>-) Possible Positions



strong chemiluminescence (CL) in aqueous cellosolve. So coelenterazine disulfonate was identified as the luciferin of *W. scintillans.*<sup>26</sup> Coelenterazine disulfonate has an imidazopyrazinone skeleton (**ImPy**) (blue part in Chart 1). Besides the firefly and bacteria BL, the BL coming from **ImPy** derivatives (**ImPys**) based luciferin or luminescent substrate of photoprotein is the most known one, which covers eight phyla (radiolaria, cnidaria, chordata, ctenophora, mollusca, arthropoda, echinodermata, chaetognatha).<sup>8,27,28</sup> Analyzing the BL process is helpful to understand the corresponding bioluminescent mechanism and provides theoretical reference on biotechnology and biomedical applications.

Tsuji has proposed a possible reaction process for the BL of W. scintillans as shown in Scheme 1.<sup>11</sup> Step 1, a rapid enolization of the keto oxygen on C<sub>3</sub> of W. scintillans luciferin (I) takes place accompanying with the departure of H from N<sub>7</sub>.

Step 2a, the enol group of I is adenylated by ATP under the catalysis of luciferase (Mg<sup>2+</sup> as the cofactor) to form adenyl coelenterazine disulfonate (II). Step 2b, the AMP group is removed from II along with the formation of the anionic form of I ( $I^-$ ). Step 3, the oxygenation of  $I^-$  generates a dioxetanone intermediate (III). Step 4, due to low stability, III decomposes to coelenteramide disulfonate (IV) and  $CO_2$  along with the emission of blue light. But until now, there has not been any strong theoretical or experimental evidence to support Tsuji's proposal. The four main problems are (1) whether ADP or AMP is produced in the adenylation of  $I_{i}^{29}$  (2) the detailed mechanisms for oxygenation of  $I^-$  by  ${}^3O_2$  (step 3 shown in red box) is uncovered; (3) the decomposition process of III (step 4) is unclear; (4) the exact chemical form of light emitter is not identified. The last two steps in the proposed bioluminescent process are responsible for producing light emitter. Comprehension to the oxygenation and decomposition in steps 3 and 4, respectively, is the key to understand the *W. scintillans* BL.

The oxygenation process is not only crucial to W. scintillans BL, but also a general step in almost all oxygen-dependent BL systems, such as firefly,<sup>30</sup> sea firefly,<sup>31</sup> jellyfish,<sup>32</sup> Obelia,<sup>33</sup> and bacteria.<sup>34</sup> For the oxygenation of I<sup>-</sup>, as summarized in Scheme 2, the oxygen molecule may attack the **ImPy** skeleton in I<sup>-</sup> at three possible positions,  $C_2^{35-41} C_3^{42,43}$  or  $C_5^{44,45}$  to form 2-, 3-, or 5-peroxide anion adduct, respectively. Which is the most likely position? Besides, there are two possible pathways when  ${}^{3}O_{2}$  attacks I<sup>-</sup>, i.e., a direct attacking way without single electron transfer (SET) and a radical-based pathway with SET from  $I^-$  to  ${}^3O_2$  via  $O_2^{\bullet-,46}$  Which one is more feasible?  $I^-$  is a derivative of ImPy as shown in Chart 1. Although oxygenation of I<sup>-</sup> itself has not been investigated, the possible positions of oxygenation of some other ImPys were previously studied by several groups.<sup>41,46-51</sup> Branchini et al. detected the presence of superoxide anion in the chemical model reaction of firefly BL, and pointed out that firefly BL is initiated by a SET from firefly luciferin to molecule oxygen.<sup>52</sup> The SET-oxygenation mechanism has been proposed for the coelenterazine BL reaction.<sup>49,53</sup>

The core structure in III is the four-membered ring (1,2dioxetanone), which is important not only for W. scintillans BL, but for the BL of firefly,<sup>30</sup> seafirefly,<sup>31</sup> jellyfish,<sup>32,54</sup> *Obelia*,<sup>50,53,55,56</sup> earthworm,<sup>6</sup> *Latia*,<sup>57</sup> and the CL of luminol,<sup>58</sup> AMPPD,<sup>59</sup> acridinium ester,<sup>60</sup> and so on. The derivatives of 1.2-dioxetanone (dioxetanones) have been researched extensively and assumed to decompose with two kinds of forms, neutral and anionic. Our previous studies have concluded for several dioxetanes<sup>61</sup> or dioxetanones.<sup>30,31,59,62-64</sup> The thermolysis of their neutral form follows the two-step biradical (entropic trapping) mechanism,  $^{30,31,61-64}$  where the C–C' bond cleaves after the complete breakage of O-O' bond. The decomposition process involves two transition states (TSs) and one biradical intermediate. In contrast to the negligible charge transfer (CT) in the decomposition of the neutral dioxetanones/dioxetanes, the decomposition of the anionic ones is induced by an obvious CT. It abides by the gradually reversible charge-transfer-induced luminescence (GRCTIL) mechanism.<sup>30</sup> In the GRCTIL mechanism, the CT and back CT (BCT) occur gradually concerted with the ruptures of O-O'and C-C' bonds, respectively. The "double crossing" between the potential energy surfaces (PESs) of the ground  $(S_0)$  and the first singlet  $(S_1)$  states is responsible for producing the light emitter.

Herein, the detailed processes and mechanisms of the I oxygenation and the III decomposition were thoroughly investigated by means of reliable density functional theory (DFT) and time-dependent DFT (TDDFT). The decomposition of III was also simulated by nonadiabatic molecular dynamics approach and the quantum yield of chemiexcitation was estimated. The knowledge obtained from *W. scintillans* BL in the current study could be expanded to a large area, because the core structure, **ImPy** skeleton in  $I^-$  is universal in the luciferin of about eight phyla of bioluminescent organisms.

# 2. COMPUTATIONAL DETAILS

Although the 3- and 5-peroxide intermediates were not detected in experimental investigations for the oxygenation of ImPys,<sup>47</sup> the C<sub>3</sub>-oxygenation could not be ruled out rashly. Because the dioxetanone intermediate can be formed directly in one-step reaction instead of through a 3-peroxide intermediate, but not be directly produced by C<sub>5</sub>-oxygenation without any intermediates (Scheme 2). The detailed

processes of  $C_2$ - and  $C_3$ -oxygenation of I were considered in this study. The reactant model is the trianion state of I ( $I^{3-}$ ), whose amide on  $N_7$ and two sulfonic group were all deprotonated, since the benzenesulfonic acid is a strong one. (The oxygenation of the monoanion state of  $I(I^-)$  with protonated sulfonic group  $(-OSO_3H)$ were also computed and the mechanism showed no obvious differences from the one of  $I^{3-}$ . (see Supporting Information (SI)) The equilibrium geometries of reactant complex (RC) formed by the I and superoxide anion radicals as starting reacting point TS, intermediate (Int), and product on the So-state surface were fully optimized with the Coulomb-attenuated hybrid exchange-correlation functional (CAM-B3LYP)<sup>65</sup> in DMSO with polarizable continuum model (PCM). For currently targeted systems, adding dispersion correction to CAM-B3LYP does not much change the calculated results (see SI). The unrestricted open-shell approximation was used to treat triplet- and singlet-state biradical species, in which the brokensymmetry technique was introduced to treat singlet-state biradical. The calculations of harmonic frequency, zero-point vibrational energy (ZVE), and thermodynamics data were carried out at the same theoretical level as the corresponding optimization. The intrinsic reaction coordinates (IRC) was performed to ensure that each TS connects the correct reactant and product. Mulliken population analysis was carried out to observe the variations of charge populations along the reaction path. The ET rate  $(k_{\rm ET})$  was estimated based on Marcus theory<sup>66</sup> (the computational method is detailedly described in SI).

For the decomposition of III, the protonated (III<sup>2-</sup>-H) and deprotonated (III<sup>3-</sup>) forms of amide on  $N_1$  atom with two - OSO<sub>3</sub><sup>-</sup> were used as the reactant models. All the electronic structure computations including geometric optimizations, vibrational frequency analyses for stationary points, and IRC on the S<sub>0</sub>-state surface were carried out in gas phase at the same theoretical levels as did in the oxygenation of  $I^{3-}$ . The Mulliken population analysis was performed to identify any possible CT characteristics. The linear response timedependent density functional theory (LR-TDDFT) was used to obtain the energies and transition properties of the S1 state based on the optimized geometric structure on the So-state PES. This method has been proved to be able to qualitatively describe the decomposition of dioxetanones.<sup>30,31,62</sup> For comparison, the decomposition processes of another two forms (III and  $III^-$  with  $-OSO_3H^{\hat{}}$ ) of the  $\hat{W}$ . scintillans dioxetanone were also investigated by the same methods (see SI for detailed computational results). The chemiluminescent decomposition of III was also simulated by the on-the-fly trajectory surface hopping (TSH) dynamics in the framework of Tully's fewest approach. In the TSH simulation, SF-DFT<sup>67-69</sup> was used to calculate the potential energies of the S<sub>0</sub> and S<sub>1</sub> states, and nuclear gradients, which incorporates the spin-flip (SF) scheme and the Tamm/Dancoff approximation  $(TDA)^{70-72}$  to TDDFT. It performs well in describing open-shell systems than the conventional Kohn-Sham DFT and giving the correct topology of  $S_0/S_1$  CI.<sup>73,74</sup> The BHHLYP functional was selected in the SF-DFT calculation, the results of which have been compared with that of (TD) CAM-B3LYP and the second-order multiconfiguration perturbation theory (CASPT2) levels on firefly dioxetanone anion.<sup>75</sup> The nonadiabatic couplings between different electronic states were derived by numerically differentials of the wave function overlaps between two successive time steps with a small duration.<sup>76,77</sup> The Newton equations were integrated using the velocity Verlet algorithm<sup>78-80</sup> with a constant time step of 0.5 fs and a maximum simulation time of 100 fs. In the initial samplings, random velocities at a constant temperature of 300 K and at a fixed TS geometry were sampled as initial condition, according to the algorithm proposed by Sellner et al.<sup>81</sup> Thus, the statistical behaviors of trajectories can be described by the canonical ensemble (NVT). The ZVE of TS (11.5 eV of  $III^{3-}$ , 12.1 eV for  $III^{-}$ ) was chosen as the initial kinetic energy. The above nonadiabatic molecular dynamics approach was also adopted to estimate the quantum yield of the light emitter (the quantum yield of chemiexcitation,  $\Phi_{\rm S}$ ) in the thermolysis of III<sup>3–</sup> and III-. This approach has been successfully employed to the chemiluminescent decomposition of firefly dioxetanone anion.<sup>30</sup>  $\Phi_s$  is also estimated based on Landau-Zener formula (see details in the

SI).<sup>82–84</sup> The solvent effect of III<sup>2–</sup> and III<sup>3–</sup> in DMSO were considered with PCM model. The 6-31G(d,p) basis set<sup>85,86</sup> was employed in all of the electronic structure computations and dynamics simulations, since our previous studies have indicated that the larger basis sets (6-31+G(d,p) and 6-31++G(d,p)), do not improve the calculated results.<sup>30</sup>

The DFT and LR-TDDFT computations were performed using the Gaussian 09<sup>87</sup> program suite. The TSH simulations were performed by Newton-X program package<sup>88,89</sup> with the electronic energies and nuclear gradients computed by the SF-DFT method in GAMESS-US.<sup>90</sup> The interface connects Newton-X and GAMESS for TSH simulation was added and the modification for CIOVERLAP module was done by previous work.<sup>75</sup>

## 3. RESULTS AND DISCUSSION

3.1. The Oxygenation of I<sup>3-</sup>. As mentioned above, the SET-oxygenation has been considered to be a general process in the oxygen-dependent BL system, and occurs easily for the dianion state of ImPys,<sup>46</sup> so oxygenation of I<sup>3-</sup> may be initiated by a SET process from  $I^{3-}$  to  ${}^{3}O_{2}$  to produce the superoxide anion  $O_2^{\bullet-}$ . The change of Gibbs free energy ( $\Delta G_{\text{SET}}$ ) of the SET process between  $I^{3-}$  and  ${}^{3}O_{2}$  was estimated in gas phase and four solvents with different dielectric coefficients and summarized in Table S1. It can be found that the SET reaction is endergonic in gas phase ( $\Delta G_{\text{SET}} = 22.2 \text{ kcal mol}^{-1}$ ), but mightily exergonic in polar solvents ( $\Delta G_{\text{SET}} = -2.8, -21.6,$ -21.7, and -22.1 kcal mol<sup>-1</sup> in benzene, CH<sub>3</sub>CN, DMSO, and water, respectively). This implies that the SET process is strongly influenced by the surrounding reaction field and probably occurs in a polar solvent thermodynamically. The estimated  $k_{\rm ET}$  from  ${\rm I}^{3-}$  to  ${}^3{\rm O}_2$  based on Marcus theory were also summarized in Table S1. The estimated value of  $k_{\rm ET}$  (6.29  $\times 10^8$  s<sup>-1</sup>) gives an evidence that SET can take place easily in DMSO, which is a common polar solvent used in experimental studies of ImPys BL.<sup>26,28</sup> So in the following discussion, the oxygenation of  $I^{3-}$  initiated by SET was investigated in DMSO. The SET between  ${}^{3}O_{2}$  and  $I^{3-}$  in DMSO generates  ${}^{3}[I^{\bullet 2-\cdots}]$ 

 $O_2^{\bullet-}]^{.91-94}$  The  ${}^{3}[I^{\bullet2-}\dotsO_2^{\bullet-}]$  must undergo a spin inversion to produce singlet dioxetanone III that is considered to decompose with light emission. Generally, the spin-alignment of  ${}^{3}[I^{\bullet2-}\dotsO_2^{\bullet-}]$  occurs via two pathways (Scheme 3). On one

# Scheme 3. Two Pathways for the Spin-Alignment of the Exciplex Forming the Singlet-State Dioxetanone Complex III<sup>3-</sup>

$$\mathbf{I}^{3-} + {}^{3}\mathbf{O}_{2} \underbrace{\overset{\mathbf{SET}}{\mathsf{DMSO}}}^{3} \left[ \mathbf{I}^{2-} \cdots \mathbf{O}_{2}^{--} \right] \underbrace{\overset{\text{weak electronic coupling}}{\underset{\mathbf{I}^{2-} + \mathbf{O}_{2}^{--}}}^{1} \left[ \mathbf{I}^{2-} \cdots \mathbf{O}_{2}^{--} \right] \xrightarrow{\mathbf{III}^{3-}}_{\text{diffused radicals}}$$

hand, the weak electronic coupling may result in an extremely small energy separation or almost degenerate between  ${}^{3}[\mathbf{I}^{\bullet 2-}\cdots O_{2}^{\bullet -}]$  and  ${}^{1}[\mathbf{I}^{\bullet 2-}\cdots O_{2}^{\bullet -}]$ , so that the hyperfine interaction between the electron spins and nuclear spins may cause the mixing of  ${}^{3}[\mathbf{I}^{\bullet 2-}\cdots O_{2}^{\bullet -}]$  and  ${}^{1}[\mathbf{I}^{\bullet 2-}\cdots O_{2}^{\bullet -}]$ , when the electron exchange interaction is smaller than the hyperfine coupling.<sup>95,96</sup> On the other hand,  ${}^{3}[\mathbf{I}^{\bullet 2-}\cdots O_{2}^{\bullet -}]$  is usually diffused efficiently by bulk solution, then the two free doublet radicals ( $\mathbf{I}^{\bullet 2-}$  and  $O_{2}^{\bullet -}$ ) emerge.  ${}^{1}[\mathbf{I}^{\bullet 2-}\cdots O_{2}^{\bullet -}]$  may be produced via the recombination of  $\mathbf{I}^{\bullet 2-}$  and  $O_{2}^{\bullet -}$  including spin inversion. As Isobe et al. pointed out, "Provided that there is no spin catalysis by paramagnetic additives such as transition-metal ions, the recombination of freely diffusing radicals is the only pathway for the spin-symmetry-allowed singlet generation from the triplet pair."<sup>46</sup> It is obvious that <sup>1</sup>[ $I^{\bullet 2-}...O_{2}^{\bullet-}$ ] is crucial for the SET-oxygenation of  $I^{3-}$ . The S<sub>0</sub> potential energy curves (PECs) of SET-oxygenations of  $I^{3-}$  at C<sub>2</sub>- and C<sub>3</sub>-position using <sup>1</sup>[ $I^{\bullet 2-}...O_{2}^{\bullet-}$ ] as initial state are shown in Figures 1 and 4. The changes of key geometric parameters and the Mulliken charge population on the  $O_{2}^{\bullet-}$  and  $I^{\bullet 2-}$  moieties along the two reaction paths were summarized in Figure 2. The electronic states were described with natural orbitals (NOs) and summarized in Tables S2 and S3. The energies of the first triplet excited state ( $T_1$ ) state ( ${}^{3}[I^{\bullet 2-}...O_{2}^{\bullet-}]$  as initial state) based on the points of S<sub>0</sub> PES were displayed in Figures 3 and 5. The PECs of C<sub>2</sub>- and C<sub>3</sub>-oxygenation of I<sup>-</sup> were displayed in Figures S3 and S4.

3.1.1. The C<sub>2</sub>-Oxygenation of  $I^{3-}$ . <sup>1</sup>[ $I^{\bullet 2-} \cdots O_{2}^{\bullet -}$ ] as Initial State. As shown in Figure 1, the singlet exciples,  $[I^{\bullet 2-} \cdots O_2^{\bullet-}]$ (RC), is formed by electrostatic force and van der Waals interaction on the S<sub>0</sub> PES of C<sub>2</sub>-oxygenation. The  $O_a - O_b$  bond distance in the RC is 1.236 Å and longer than it in  ${}^{3}O_{2}$  (1.205 Å), which implies  $O_2^{\bullet-}$  rather than  ${}^3O_2$  attacks the substrate  $I^{3-}$ . There is a partial negative charge transferred from  $O_2^{\bullet-}$  to  $I^{\bullet 2^-}$  during the formation of RC, because the Mulliken charge on the  $O_2^{\bullet-}$  moiety,  $\rho(O_2^{\bullet-})$ , changes from -1lel in  $O_2^{\bullet-}$  to -0.32lel in RC. From RC to the final adduct III<sup>3-</sup>, four TSs and three Ints were located. Vibrational analysis indicates that TS1(214.6*i*), TS3(151.5*i*), and TS4(155.0*i*) correspond to the stretching of  $C_2$ - $O_a$ ,  $C_3$ - $O_b$ , and  $C_3$ - $N_4$  bond, respectively (see Figure 1). But the imaginary vibrational modes of TS2 (2335.5i) could not be obviously attributed to a single imaginary canonical vibrational mode, both C2-Oa and Oa-O<sub>b</sub> stretch synchronously in TS2.

As shown in Figure 2A, the C<sub>2</sub>-oxygenation reaction starts with the gradual shortening of  $C_2-O_a$  bond along with the torsion of  $C_3-C_2-O_a-O_b$  toward to the direction of the closure of four-membered ring. After TS2, the variation curves of  $C_2 - O_a$  remains almost flat from 4.308 amu<sup>1/2</sup> bohr until the formation of final product  $III^{3-}$ . In the flat region of  $C_2-O_a$ bond curve, the  $C_3 - O_b$  bond and the dihedral  $C_3 - C_2 - O_a - O_b$ keep almost unchanged at the beginning (between 4.308 and 38.010 amu<sup>1/2</sup> bohr), then changes quickly to form the fourmembered ring. Int2 (with 1.399 Å  $C_2-O_3$ , -74.2 ° $C_3-C_2 O_a - O_b$ , and  $-1.00 lel \rho(O_2^{\bullet-})$ ) was located in this platform, which is the 2-peroxide anion. The 2-peroxide was also trapped in previous experimental investigation on the oxygenation of ImPys.<sup>47</sup> In the vicinity of Int3, the variation curves of  $C_3 - O_b$ and  $C_3 - C_2 - O_a - O_b$  remain flat again, which indicates that the four-membered ring is formed. The C3-N4 bond keeps unchanged from the beginning of C2-oxygention reaction, and starts to break to produce III<sup>3-</sup> after Int3 accompanying with the minor adjustments of C2-Oa and C3-Ob bonds. So the cleavage of C<sub>3</sub>-N<sub>4</sub> bond and the closure of four-membered ring are stepwise. The expectation value of  $S^2$  operator ( $\langle S^2 \rangle$  in Figure 1) demonstrates that the  $S_0$  state has obvious biradical characteristic from RC to TS2 (1.02 for RC, 1.00 for TS1, 0.99 for Int1, and 0.82 for TS2). The PEC between TS1 and TS2 is very short and almost flat. After TS2, the biradical character of molecule disappears, and the S<sub>0</sub> energy suddenly decreases to form Int2. Among all the energy barriers on the S<sub>0</sub> PEC, the one from RC to Int1 is the highest (9.8 kcal  $mol^{-1}$ ). So this step is the rate-determining one, whose barrier is 1.4 kcal mol<sup>-1</sup> lower than that of the  $C_2$ -oxygenation of I<sup>-</sup> (shown in Figure S3).

The NOs analysis along the IRC (Table S2) have shown that the NOs are mainly located at the  $O_2^{\bullet-}$  moiety in the vicinity of



**Figure 1.**  $S_0$  PEC ( $^1[I^{\bullet 2} \dots O_2^{\bullet -}]$  as initial state) of  $C_2$ -oxygenation of  $I^{3-}$  takes place between  $I^{\bullet 2-}$  and  $O_2^{\bullet -}$  in DMSO computed at the UCAM-B3LYP/6-31G(d,p) level. The key atoms ( $N_1$ ,  $C_2$ ,  $C_3$ ,  $N_4$ ,  $O_a$ , and  $O_b$ ) are highlighted in yellow, and the dihedral angle ( $C_3 - C_2 - O_a - O_b$ ) in pink. Here and hereunder, gray, blue, white, and red balls represent carbon, nitrogen, hydrogen, and oxygen atoms, respectively. The units are angstrom (Å) for bond distance, degree (deg) for the dihedral angle, and kcal mol<sup>-1</sup> for Gibbs energy change ( $\Delta G$ ).

RC. So the  $S_0$  state exhibits  ${}^2\Pi_g$  configuration. With the approaching of  $O_2^{\bullet-}$ , the degenerate  $\pi^*$  orbital of  $O_2^{\bullet-}$  couples with the  $\pi$  orbital distributed in ImPy skeleton of I<sup>•2-</sup>, which brings about the S<sub>0</sub> state of  $(\pi, \pi^*)$  configuration. Then the S<sub>0</sub> state turns into the closed-shell configuration after TS2. Therefore,  $S_0$  is denoted as  ${}^{1}({}^{2}\Pi_{\sigma}-O_{2}^{\bullet-}+{}^{2}I^{\bullet2-})/{}^{1}(\pi,\pi^{*})/$ closed-shell state. According to the variation of Mulliken population along the IRC (Figure 2C), a gradual CT from  $I^{\bullet 2-}$ to  $O_2^{\bullet-}$  takes place before TS2, and then followed by a sharply decreasing of the charge of  $\mathbf{I}^{\bullet 2^-}$ ,  $\rho(\mathbf{I}^{\bullet 2^-})$  (-2.55lel  $\rightarrow$  -2.08lel) and increasing of  $\rho(O_2^{\bullet-})$  (-0.45lel  $\rightarrow$  -0.92lel). Afterward, the Mulliken charges showed no obvious variation between 4.308 and 38.010 amu<sup>1/2</sup> bohr, where the  $\rho(O_2^{\bullet-})$  remains nearly -1|e|. It fully proves the existence of 2-peroxide anion in conjugation with the key geometric parameters (Figure 2A). Then a gradual BCT process from  $O_2^{\bullet-}$  to  $I^{\bullet 2-}$  takes place from TS3 to the final product III<sup>3-</sup>. The  $\rho(O_2^{\bullet-})$  and  $\rho(I^{\bullet 2-})$  located at the product  $III^{3-}$  is -0.64lel and -2.36lel (Figure 1). The C<sub>2</sub>oxygenation pathway is consistent with experimental observation of the two complex signals (2-peroxide and dioxetanone intermediates) by <sup>13</sup>C NMR spectroscopy.<sup>47</sup> It is noteworthy that the distance of  $C_2 - O_a$  bond in Int1 is 1.558 Å and in III<sup>3-</sup> is 1.547 Å (Figure 1), which indicates that the  $C_2-O_a$  bond is already formed in Int1. But Int1 has biradical character ( $\langle S^2 \rangle$  = 0.99), and then produces to the 2-peroxide anion (Int2) via TS2, which implies that the 2-peroxide anion is formed via a

nucleophilic addition followed by TS2 with biradical annihilation rather than the generally proposed radical coupling in coelenterate, <sup>53</sup> *Cypridina*.<sup>49</sup>

 ${}^{3}[I^{\bullet 2-}\cdots O_{2}^{\bullet -}]$  as Initial State. As shown in Figure 3, at beginning, the  $T_1$  state takes on  ${}^2\Pi_{\sigma}$  configuration, then changes to  ${}^3(n, \pi^*)$ . Here, n is the p orbitals of  $O_a$  and  $O_b$ atoms, and  $\pi^*$  orbital mainly represents the  $\pi$ -antibonding orbital of ImPy ring in  $I^{\bullet 2-}$ . Finally, the  $T_1$  state turns into  ${}^3(\pi,$  $\pi^*$ ) configuration with the formation of four-membered ring. The  $\pi$  and  $\pi^*$  distribute in **ImPy** ring. So T<sub>1</sub> is represented as  ${}^{3}({}^{2}\Pi_{g}$ - $\Omega_{2}^{\bullet-} + {}^{2}I^{\bullet2-})/{}^{3}(n, \pi^{*})/{}^{3}(\pi, \pi^{*})$ . The energies in the S<sub>0</sub> and T1 states are almost degenerate in the short and flat biradical region (between TS1 and TS2), and the two PECs begin to separate after TS2, which implies a S/T surface intersection nearby TS2. Unfortunately, the currently employed method does not have the capability to optimize the geometry of S/T crossing point, and the multiconfiguration method is not available to handle such a large system. According to Figure 3, <sup>1</sup>III<sup>3-</sup> may be produced by <sup>3</sup>[ $I^{\bullet 2-} \cdots O_2^{\bullet -}$ ] via two pathways. In the first way,  ${}^{3}\text{III}{}^{3-}$  is produced from  ${}^{3}[I^{\bullet 2-}\cdots O_{2}^{\bullet-}]$  by overcoming a barrier of 59.2 kcal mol<sup>-1</sup>. Then, <sup>1</sup>III<sup>3-</sup> is formed after <sup>3</sup>III<sup>3-</sup> undergoes a ISC process. In the second way shown by cyan arrow in Figure 3, <sup>1</sup>III<sup>3-</sup> is formed directly by <sup>3</sup>[I<sup> $\cdot$ 2</sup>...  $O_2^{\bullet-}$  via the ISC at the S/T crossing point, which needs to overcome a barrier of 12.2 kcal mol<sup>-1</sup>. The barrier of the second way is much lower than that of the first way. Moreover,



Figure 2. Key geometric parameters and the Mulliken charges on the  $O_2^{\bullet-}$  and  $I^{\bullet 2-}$  moieties in the  $C_2$ -oxygenation (A, C) and  $C_3$ -oxygenation (B, D) pathways along the IRC calculated in DMSO at the UCAM-B3LYP/6-31G(d,p) level.

on the basis of El-Sayed's rule, the rate of ISC is relatively large when the transition involves a change of molecular orbital type.<sup>97</sup> A change of the molecular orbital type indeed exists in the second way ( ${}^{3}(n, \pi^{*})$  and  ${}^{1}(\pi, \pi^{*})$  in the degenerate region around TS2) as shown in Figure 3. So  ${}^{1}\text{III}{}^{3-}$  prefers to be formed from  ${}^{3}[\mathbf{I}^{\bullet 2-}\cdots O_{2}^{\bullet -}]$  via the second way.

3.1.2. The  $C_3$ -Oxygenation of  $I^{3-}$ .  ${}^{1}[I^{\bullet 2-}...O_{2}^{\bullet -}]$  as Initial State. As shown in Figure 4, the singlet exciplex,  ${}^{1}[I^{\bullet 2-}...O_{2}^{\bullet -}]$  (RC), is located in  $C_3$ -oxygenation pathway with the  $C_3-O_b$  bond of 2.792 Å. The distance of  $O_a-O_b$  (1.223 Å) and the negative charges -0.15lel indicate that  $I^{3-}$  is oxidized by  $O_2^{\bullet -}$  instead of  ${}^{3}O_2$ , and that partial negative charges are transferred from  $O_2^{\bullet -}$  to  $I^{\bullet 2-}$  with the approaching of  $O_2^{\bullet -}$ , which is similar to the  $C_2$ -oxygenation pathway. However, from singlet RC, only one TS (1100.5*i*) is located, which corresponds to the simultaneous stretching of  $C_3-O_b$  and  $C_3-N_4$  bond. There is no peroxide intermediate formed, which is different from the  $C_2$ -oxygenation pathway.

On the basis of the variations of dominant geometric parameters in Figure 2B, the C<sub>3</sub>-oxygenation reaction starts with the formation of  $C_3-O_b$  along with the torsion of  $C_2-C_3-O_b-O_a$  toward to the direction of the closure of the four-membered ring. After TS, a platform appears on the variation curve of  $C_3-O_b$  from 1.519 amu<sup>1/2</sup> bohr to the formation of final product, which is similar to the situation of the C<sub>2</sub>-

oxygenation. Different to the C<sub>2</sub>-oxygenation pathway, in this flat region, the C<sub>2</sub>–O<sub>a</sub> bond is getting shorter and C<sub>3</sub>–N<sub>4</sub> getting longer to produce the final product III<sup>3–</sup> directly. So the 3-peroxide anion does not exist in this process. This also indicates that the cleavage of C<sub>3</sub>–N<sub>4</sub> bond and the closure of four-membered ring are concerted. The  $\langle S^2 \rangle$  (Figure 4) demonstrates that S<sub>0</sub> state has obvious biradical characteristics from RC (0.94) to TS (0.55), but the biradical character of molecule disappears after TS. The activation barrier of C<sub>3</sub>-oxygenation is 18.4 kcal mol<sup>-1</sup>, which is 0.8 kcal mol<sup>-1</sup> lower than that of the C<sub>3</sub>-oxygenation of I<sup>-</sup> (see Figure S4), but 8.6 kcal mol<sup>-1</sup> higher than that of the C<sub>2</sub>-oxygenation of I<sup>3–</sup>. So from the viewpoint of thermodynamics, the C<sub>2</sub>-oxygenation of I<sup>3–</sup> is energetically feasible using <sup>1</sup>[I<sup>+2–</sup>…O<sub>2</sub><sup>•–</sup>] as initial state.

Similar to the case of C<sub>2</sub>-oxygenation, the S<sub>0</sub> state behaves a <sup>2</sup>Π<sub>g</sub> configuration, then gradually presents ( $\pi$ ,  $\pi^*$ ) configuration along with the mixing between  $\pi^*$  orbital of O<sub>2</sub><sup>•-</sup> and  $\pi$  orbital of the **ImPy** ring in **I**<sup>•2-</sup>. After TS, the S<sub>0</sub> state turns into the closed-shell configuration in which all of the bonding and nonbonding orbitals are doubly occupied. So the S<sub>0</sub> state is denoted as  ${}^{1}({}^{2}\Pi_{g}$ -O<sub>2</sub><sup>•-</sup> +  ${}^{2}$ **I**<sup>•2-</sup>)/ ${}^{1}(\pi, \pi^*)$ /closed-shell state. Correspondingly, the charge gradually transfers from **I**<sup>•2-</sup> to O<sub>2</sub><sup>•-</sup> moieties in the vicinity of RC accompanied by the formation of the first C–O bond (i.e., C<sub>3</sub>–O<sub>b</sub> bond) as shown in Figure 2D. After TS2,  $\rho(O_2^{•-})$  sharply increases from -0.66l



Figure 3.  $S_0$  and  $T_1$  PECs of the C<sub>2</sub>-position oxygenation of  $I^{3-}$  in DMSO calculated at the UCAM-B3LYP/6-31G(d,p) level.



**Figure 4.**  $S_0$  PEC ( ${}^3[I^{\bullet 2} \dots O_2^{\bullet -}]$  as initial state) of  $C_3$ -oxygenation of  $I^{3-}$  occurs between  $I^{\bullet 2-}$  and  $O_2^{\bullet -}$  computed in DMSO at the UCAM-B3LYP/ 6-31G(d,p) level. The key atoms ( $N_1$ ,  $C_2$ ,  $C_3$ ,  $N_4$ ,  $O_a$  and  $O_b$ ) are highlighted in yellow, and the dihedral angle ( $C_2-C_3-O_b-O_a$ ) in pink.

el to -0.96lel, and correspondingly,  $\rho(\mathbf{I}^{\bullet 2-})$  decreases from -2.34lel to -2.04lel. Then there is a gradual BCT from  $O_2^{\bullet-}$  to  $\mathbf{I}^{\bullet 2-}$  coupled with the formation of the four-membered ring.

<sup>3</sup>[ $I^{\bullet 2-\cdots}O_2^{\bullet -}$ ] as Initial State. According to the NOs analysis, the T<sub>1</sub> state (Figure 5) changes from the <sup>2</sup> $\Pi_g$  configuration (near the RC) to the (n,  $\pi^*$ ) configuration, and takes on a little ( $\pi$ ,  $\pi^*$ ) configuration neighboring the rupture of C<sub>3</sub>-N<sub>4</sub> bonds from 35.508 to 36.304 amu<sup>1/2</sup> bohr. So the T<sub>1</sub> is denoted as <sup>3</sup>(<sup>2</sup> $\Pi_g$ -O<sub>2</sub><sup>•-</sup>+ <sup>2</sup>I<sup>•2-</sup>)/<sup>3</sup>(n,  $\pi^*$ )/<sup>3</sup>( $\pi$ ,  $\pi^*$ ) state. Similar to the C<sub>2</sub>oxygenation, the T<sub>1</sub> state is almost degenerate with the S<sub>0</sub> state near TS. After TS, the energies of the T<sub>1</sub> state are apparently higher than those of the S<sub>0</sub> state. This also implies a S/T surface intersection around TS.  ${}^{3}[\mathbf{I}^{\bullet 2-}\cdots O_{2}^{\bullet -}]$  can generate  ${}^{1}\mathbf{III}{}^{3-}$  via two pathways. (1)  ${}^{3}[\mathbf{I}^{\bullet 2-}\cdots O_{2}^{\bullet -}]$  needs to pass over a barrier of 71.8 kcal mol<sup>-1</sup> to form  ${}^{3}\mathbf{III}{}^{3-}$ , then  ${}^{3}\mathbf{III}{}^{3-}$  undergoes a ISC to produce  ${}^{1}\mathbf{III}{}^{3-}$  (Figure 5). (2)  ${}^{3}[\mathbf{I}^{\bullet 2-}\cdots O_{2}^{\bullet -}]$  overcomes a barrier of 26.3 kcal mol<sup>-1</sup> then goes through a ISC near TS to form  ${}^{1}\mathbf{III}{}^{3-}$  (the cyan arrow in Figure 5). The latter pathway has much lower barrier. And the transition between  ${}^{1}(\pi, \pi^{*})$ and  ${}^{3}(n, \pi^{*})$  is much likely to occur in the vicinity of TS based on the El-sayed's rule. Hence,  ${}^{1}\mathbf{III}{}^{3-}$  is probably produced via ISC through this S/T crossing point near TS from  ${}^{3}[\mathbf{I}^{\bullet 2-}\cdots O_{2}^{\bullet -}]$ .



Figure 5.  $S_0$  and  $T_1$  PECs of the C<sub>3</sub>-position oxygenation of  $I^{3-}$  in DMSO calculated at the UCAM-B3LYP/6-31G(d,p) level.



**Figure 6.** (A) The S<sub>0</sub> PECs for III<sup>2–</sup> decomposition; (B) The approximate diabatic PECs of the  ${}^{1}(\sigma, \sigma^{*})$ ,  ${}^{1}(\pi, \pi^{*})$ , and  ${}^{1}(n, \sigma^{*})$  states with the adiabatic PECs of the S<sub>0</sub> and S<sub>1</sub> states (inset); (C) The variation of key geometric parameters; (D) The populations of Mulliken charges on CO<sub>2</sub> and IV<sup>2–</sup> along the IRC computed in gas phase at the UCAM-B3LYP/6-31G(d,p) level. The key atoms (C<sub>2</sub>, C<sub>3</sub>, O<sub>a</sub>, and O<sub>b</sub>) are highlighted in yellow, and the dihedral angle (O<sub>b</sub>-C<sub>3</sub>-C<sub>2</sub>-O<sub>a</sub>) in pink.

3.1.3.  $C_2$ -Oxygenation vs  $C_3$ -Oxygenation. As discussed above, in formation process of  ${}^{1}\text{III}{}^{3-}$  from  ${}^{1}[\text{I}^{\bullet 2-}\cdots\text{O}_{2}^{\bullet-}]$ , the 2peroxide anion is formed, so the  $C_2$ -oxygenation of  $I^{3-}$  is stepwise. The C<sub>3</sub>-oxygenation is concerted because there is not any peroxide anion formed. This can also be found by comparing the changes of the C3-N4 bond and the closure of the four-membered ring. The C2-oxygenation is energetically feasible and agrees with experimental conclusion on 2-peroxide and dioxetanone.<sup>47</sup> The assumption of radical-coupling reaction between  $O_2^{\bullet-}$  and  $I^{\bullet 2-}$  (or luminescent substrate radical)<sup>49,53</sup> is not supported by current results. Int2 is directly transformed from Int1 via nucleophilic addition followed by TS2 of biradical annihilation rather than by the coupling of  $O_2^{\bullet-}$  and  $I^{\bullet 2-}$ radicals in C2-oxygenation. However, the spin-orbital coupling cannot be achieved by the current theoretical method, so the formation rate of  ${}^{1}\mathbf{III}^{3-}$  from  ${}^{3}[\mathbf{I}^{\bullet 2-}\cdots \mathbf{O}_{2}^{\bullet -}]$  was not computed. Nevertheless,  ${}^{3}[I^{\bullet 2} \cdots O_{2}^{\bullet -}]$  needs about 26.3 kcal mol<sup>-1</sup> to approach the  $C_3$ -ISC, which is much higher than that to the  $C_2$ -ISC (12.2 kcal mol<sup>-1</sup>). Using <sup>3</sup>[I<sup>•2-...</sup>O<sub>2</sub><sup>•-</sup>] as the initial RC, the C<sub>2</sub>-oxygenation state is energetically more feasible than C3-oxygenation. But we do not have strong evidence to judge whether  ${}^{1}[\mathbf{I}^{\bullet 2-}\cdots \mathbf{O}_{2}^{\bullet -}]$  or  ${}^{3}[\mathbf{I}^{\bullet 2-}\cdots \mathbf{O}_{2}^{\bullet -}]$  is more favorable in C<sub>2</sub>-oxygenation based on the current calculations.

3.2. The Chemiluminescent Decomposition of W. scintillans Dioxetanone III. The thermolysis and chemiexcitation mechanisms of both III<sup>2-</sup> and III<sup>3-</sup> were considered in gas phase and DMSO. The calculated PECs of S<sub>0</sub> and S<sub>1</sub> states, key geometric parameters, and the populations of Mulliken charges were depicted in Figures 6 and 7. Time evolution of the average distribution of the  $O_a-O_b$  and  $C_2-C_3$  bond distances, the Mulliken charges of  $CO_2$  and  $IV^{2-}/IV^{3-}$  moieties, and the fractional occupations of the S<sub>0</sub> and S<sub>1</sub> states during the TSH simulation were shown in Figures 8 and 9, respectively. The comparisons of the S<sub>0</sub> PECs in gas phase and DMSO for III<sup>2-</sup> and III<sup>3-</sup> decompositions were shown in Figure S6. The main NOs changes during III<sup>2-</sup> and III<sup>3-</sup> decompositions were summarized in Tables S4 and S5. For comparison, the decompositions of III and III<sup>-</sup> with  $(-OSO_3H)_2$  were further investigated and related data were summarized in Figures S5, S7, S10, and S11.

3.2.1. Thermolysis and Chemiexcitation of III<sup>2-</sup>. Two TSs and one Int were located on the  $S_0$  PES during the thermolysis of  $III^{2-}$  in gas phase (Figure 6A). The imaginary vibrational modes of TS1 (438.8i) and TS2 (1067.3i) correspond to the stretching of  $O_a - O_b$  and  $C_2 - C_3$  bond, respectively. As shown in Figure 6C, the decomposition of III<sup>2-</sup> starts with the elongation of O<sub>a</sub>-O<sub>b</sub> bond. But a platform appears on the variation curve of O<sub>a</sub>-O<sub>b</sub> bond between 2.841 and 10.903  $amu^{1/2}$  bohr. In this region, the  $O_a - O_b$  bond distance is about 2.7 Å, which indicates that the  $O_a - O_b$  bond is completely ruptured. The  $C_2-C_3$  bond keeps unchanged from the beginning of III<sup>2-</sup>decomposition (Figure 6C), then begins to elongate to form final product  $IV^{2-}$  after 4.176 amu<sup>1/2</sup> bohr. An Int  $(O_a - O_b \text{ of } 2.681 \text{ Å and } C_2 - C_3 \text{ of } 1.529 \text{ Å})$  was located in this flat stage. These results indicate that the III<sup>2-</sup> decomposition is a typical stepwise process, which is similar to the neutral decomposition of simple dioxetanones.  $^{30,31,62,65-67}$  The dihedral angles of  $O_b-C_3-C_2-O_a$  $(5.0^{\circ} \text{ in III}^{2-}, 8.2^{\circ} \text{ in TS1}, 14.7^{\circ} \text{ in Int}, 12.4^{\circ} \text{ in TS2})$  indicate that the four-membered ring keeps nearly planar during the whole process. The  $\langle S^2 \rangle$  values (0.75 for TS1, 1.00 for TS2, 1.01 for Int in Figure 6A) demonstrates that TS1, TS2, and Int all have obvious biradical characters. The S<sub>0</sub> PEC between TS1

and TS2 is nearly flat with biradical characteristic. The energy barriers from III<sup>2–</sup> to Int and the final product are 23.5 and 23.2 kcal mol<sup>-1</sup>, respectively. So the first step is the rate-determining step, which is in line with the experimental results of other dioxetanes and dioxetanones (about 20.0 kcal mol<sup>-1</sup> see SI in ref 30). As shown in Figure 6D, there is no obvious CT from IV<sup>2–</sup> to CO<sub>2</sub> (<0.1lel) moiety on the S<sub>0</sub>-state PEC. But a sudden CT/BCT about –0.30lel appears between 4.691 and 6.765 amu<sup>1/2</sup> bohr, which is different from non-CT decompositions of the simple dioxetanones.<sup>30,31,65–67</sup>

According to the adiabatic PECs (the inset in Figure 6B), the  $S_0$  and  $S_1$  curves approach close to each other as the elongation of O<sub>a</sub>-O<sub>b</sub> bond, and then becomes nearly parallel at the biradical platform region between -1.055 and 4.692 amu<sup>1/2</sup> bohr. The NOs analysis (Table S4) along the IRC indicates that S<sub>0</sub> surface gradually changes from closed-shell to  ${}^{1}(\sigma, \sigma^{*})$ configuration, while S<sub>1</sub> state changes from  ${}^{1}(\pi, \pi^{*})$  to  ${}^{1}(n, \sigma^{*})$ configuration before TS2. Here  $\sigma$  and  $\sigma^*$  correspond to the bonding and antibonding molecular orbitals of O<sub>a</sub>-O<sub>b</sub> bond, and n refers to the out-of-plane atomic p orbital of O<sub>a</sub> and inplane atomic p orbital of  $O_{\rm b}$ . The  $\pi$  and  $\pi^*$  orbital are mainly contributed by the pyrazine ring. These variations of electronic structure are similar to those neutral decompositions of previously studied dioxetanones,  $^{30,31,65-67}$  where nearly degenerated  ${}^{1}(\sigma,\sigma^{*})$  and  ${}^{1}(n,\sigma^{*})$  PECs were located during the flat biradical region.  ${}^{30,31,61,64}$  However, the interval of the S<sub>0</sub> and S<sub>1</sub> PECs is large rather than small at the flat biradical region (Figure 6B). The energetic error is caused by the UTDDFT calculation, which is not able to describe the open-shell biradical as a single configuration method. The SF-DFT calculations can dramatically reduce the interval at the flat region although it is still not an ideal way to handle the biradical systems (see Figure S2A in section 3 of SI). The problem frequently occurs to the thermolysis of dioxetanones, for examples, the dioxetanones of firefly<sup>30</sup> and *Cypridina*.<sup>31</sup> For the firefly dioxetanone, the UTDDFT calculation predicted a large  $S_0/S_1$  energy gap (more than 20 kcal mol<sup>-1</sup>) at the flat biradical region, although its computed electronic structures and shapes of the  $S_0$  and  $S_1$  PECs (Figure S11A in ref 30) agree with the ones (Figure 2A in ref 30) calculated by the multireference method, respectively. Nevertheless, the large  $S_0/S_1$  energy gap dramatically decreases after the energy correction by the multireference methods. The UTDDFT is qualitatively reliable for analyzing the thermolysis of dioxetanones. To describe the III<sup>2-</sup> decomposition with quantitative accuracy, a multireference method has to be adopted as we have done on the firefly dioxtanone.<sup>30</sup> However, it is nearly impossible for such a large system as III<sup>2-</sup>. On the basis of our previous studies on the decompositions of several dioxetanones, 30,31,61,64 the nearly degenerated  $S_0/S_1$  PECs at the flat region are expected. In the nearly barrierless region, the reaction is mainly driven by the entropic effect and follows entropic trapping mechanism, which compels the reaction molecule to spend a long time passing through the PEC. The nearly degenerated  $S_0/S_1$  PECs could provide the possibility for the nonadiabatic coupling as reported in previous studies.<sup>30</sup> Moreover, with the  $C_2 - C_3$  bond cleavage after TS2, the  $\pi$  orbital of pyrazine ring coupled with the p orbital of  $N_1$  and the  $\sigma^*$  orbital of  $O_a - O_b$  bond, which brings some  ${}^{1}(\pi, \sigma^{*})$  character into the original  ${}^{1}(\sigma, \sigma^{*})$  configuration of the S<sub>0</sub> state and the  ${}^{1}(\sigma, \sigma^{*})$  configuration of the S<sub>1</sub> state. These states mixing between 5.038 and 9.192  $\text{amu}^{1/2}$  bohr (pink box in Figure 6B) leads to a CT character on the  $S_0$ surface as observed in Figure 6D and a smaller  $S_0/S_1$  energy gap



**Figure 7.** (A) The S<sub>0</sub> PEC for III<sup>3-</sup> decomposition; (B) The approximate diabatic PECs of the  ${}^{1}(\sigma, \sigma^{*})$ ,  ${}^{1}(\pi, \pi^{*})$ , and  ${}^{1}(\pi, \sigma^{*})$  states with the adiabatic PECs of the S<sub>0</sub> and S<sub>1</sub> states (inset); (C) The changes of key geometric parameters (O<sub>a</sub>-O<sub>b</sub> and C<sub>2</sub>-C<sub>3</sub> in Å, O<sub>b</sub>-C<sub>3</sub>-C<sub>2</sub>-O<sub>a</sub> in degrees labeled in pink); (D) The populations of Mulliken charges on the CO<sub>2</sub> and IV<sup>3-</sup> moieties along the IRC computed in gas phase at the UCAM-B3LYP/6-31G(d,p) level. The key atoms (C<sub>2</sub>, C<sub>3</sub>, O<sub>a</sub>, and O<sub>b</sub>) are highlighted in yellow, and the dihedral angle (O<sub>b</sub>-C<sub>3</sub>-C<sub>2</sub>-O<sub>a</sub>) in pink.

 $(\Delta E = 3.2 \text{ kcal mol}^{-1})$  at 6.423 amu<sup>1/2</sup> bohr (Figure 6B), which is different to the neutral decompositions of simple dioxetanones.<sup>30,31,65-67</sup> This might be attributed to the two negative charges populated on the two – OSO<sub>3</sub><sup>-</sup>, which increases the redox potential of III<sup>2-</sup>. For comparison, we also investigated the chemiluminescent decomposition of III. A similar stepwise biradical process (Figure S5A) without no obvious CT was observed. As shown in Figure S5B, the S<sub>0</sub> and S<sub>1</sub> state of III are dominated by the non-CT closedshell/<sup>1</sup>( $\sigma$ , $\sigma$ \*) and <sup>1</sup>(n, $\pi$ \*)/<sup>1</sup>(n, $\sigma$ \*) configuration without obvious <sup>1</sup>( $\pi$ ,  $\sigma$ \*) configuration, which is very similar to the neutral decomposition of *Cypridina*<sup>31</sup> and firefly dioxetanone.<sup>30</sup>

To sum up, with the influence of the two  $-OSO_3^-$ , the decomposition and chemiexcitation process of  $III^{2-}$  is a little different to the neutral decompositions of simple dioxetanones. However, the thermolysis of  $III^{2-}$  can also take place via a stepwise-biradical process and its chemiexcitation process could be understood with the entropic trapping mechanism, which is similar to multireference-calculations of the neutral dioxetanones.<sup>30,31,65-67</sup> As shown in Figure S6A, the S<sub>0</sub> PEC is almost unaffected by DMSO, as there is no obvious CT in the  $III^{2-}$  decomposition.

3.2.2. Thermolysis and Chemiexcitation of  $III^{3-}$ . As shown in Figure 7A, only one TS with large imaginary frequency (1172.6*i*) was located on the S<sub>0</sub> PES at the UCAM-B3LYP/6-31G(d,p) level, which corresponds to the stretching of O<sub>a</sub>-O<sub>b</sub> bond. The decomposition of III<sup>3-</sup> starts with the elongation of

 $O_a - O_b$  peroxide bond (Figure 7C) like the III<sup>2-</sup> thermolysis. Then a short plateau exists between 5.259 and 7.678 amu<sup>1/2</sup> bohr in the variation curve of  $O_a - O_b$  bond. Different to the  $III^{2-}$  decomposition, in this region, the  $O_a - O_b$  bond distance is about 2.3 Å, which indicates that the  $O_a - O_b$  bond does not break completely. The  $C_2-C_3$  bond begins to break in the period of stagnation of the O<sub>a</sub>-O<sub>b</sub> variation curve instead of keeping unchanged in the III<sup>2-</sup> thermolysis. So the cleavages of  $O_a - O_b$  and  $C_2 - C_3$  bonds proceed in an asynchronousconcerted fashion. In addition, the dihedral angle of  $O_b - C_3 C_2$ - $O_a$  changes from 3.0° to 11.5°, finally back to 7.8°, which indicates that the four-membered ring remains nearly planar during the III<sup>3-</sup> decomposition. The  $\langle S^2 \rangle$  (0.38) and  $\rho(CO_2)$  $(III^{3-} \rightarrow TS \text{ is } -0.32 \text{lel} \rightarrow -0.17 \text{lel})$  of TS indicate that it has some biradical and CT characteristics. The activation energy over the TS is 14.5 kcal mol<sup>-1</sup>, which is 9.0 kcal mol<sup>-1</sup> lower than that in  $III^{2-}$  decomposition, 7.2 kcal mol<sup>-1</sup> lower than that in III decomposition (Figure S5A), and 1.3 kcal mol<sup>-1</sup> lower than that in III-decomposition (Figure S7A). Hence the decomposition of III<sup>3-</sup> is energetically more feasible for W. scintillans BL in terms of the thermodynamics. Moreover, the variations of Mulliken charges (Figure 7D) suggest a gradual CT of  $IV^{3-} \rightarrow CO_2$  with  $\Delta \rho = -0.39$ lel from -4.227 to 3.608 amu<sup>1/2</sup> bohr and a gradual BCT of  $CO_2 \rightarrow IV^{3-}$  with  $\Delta \rho$ = -0.53lel from 3.608 to 14.113 amu<sup>1/2</sup> bohr. In one word, the thermolysis of III<sup>3-</sup> is a planar asynchronous-concerted process with CT character, which is similar to the anionic



Figure 8. Time evolution of the average distribution of the  $O_a-O_b$  and  $C_2-C_3$  bond distances at 0, 20, 40, 60, 80, and 100 fs during the dynamics simulation of III<sup>3-</sup> (average over 100 trajectories) obtained by the TSH dynamics in gas phase at the SF-BHHLYP/6-31G(d,p) level.

decomposition of other dioxetanones like  $FDO^{-30}$  and  $CDO^{-31}$ .

According to the NOs analysis (Table S5), the  $S_0$  state changes from the closed-shell to  ${}^1(\pi, \sigma^*)$  configuration along with the gradual elongation of  $O_a - O_b$  bond, which is different from  $III^{2-}$  decomposition. Here  $\pi$  denotes the p $\pi$ -conjugated orbital of N<sub>1</sub> and the pyrazine ring, and  $\pi^*$  represents the  $\pi$ antibonding orbital of pyrazine ring. But in the initial stage, the main configuration of the S<sub>1</sub> state is  $(\pi, \pi^*)/(\pi, \sigma^*)$ , but changes to  ${}^{1}(\sigma, \sigma^{*})$  with the rupture of  $O_{a} - O_{b}$  bond. After the complete rupture of  $C_2-C_3$  bond, the  $S_1$  state changes to  ${}^1(\pi,$  $\pi^*$ ) configuration and the S<sub>0</sub> state returns to closed-shell configuration again. Here  $\pi$  and  $\pi^*$  stand for the bonding and antibonding molecular orbitals located in the benzopyrazine ring. So the  $S_0$  and  $S_1$  states can be approximately described by the closed-shell/ $(\sigma, \sigma^*)$  and  $(\pi, \pi^*)/(\pi, \sigma^*)/(\sigma, \sigma^*)$ configurations, respectively. (Figure 7B) The approximate diabatic PECs clearly shows a "double crossings" between the closed shell/ $(\sigma, \sigma^*)$  and  $(\pi, \pi^*)/(\pi, \sigma^*)/(\sigma, \sigma^*)$  surface, which provides the channels for nonadiabatic transition. The smallest energy gaps of  $S_0$  and  $S_1$  states in the vicinities of the double crossings are 17.7 and 15.0 kcal mol<sup>-1</sup> at IRC = -0.060and 4.860  $\operatorname{amu}^{1/2}$  bohr, respectively, which could be larger than the physical truth due to the spin contamination mentioned in III<sup>2-</sup> decomposition. The multiconfiguration methods can decrease the energy gap, but cannot be adopted to treat with such a big system  $(III^{3-})$ . A smaller energy gap about 0–3 kcal mol<sup>-1</sup> was found via the SF-DFT calculations (see Figure S2B in section 3 of SI). In addition, the gradual CT/BCT takes place simultaneously with the  $O_a-O_b$  and  $C_2-C_3$  bond cleavages, respectively. Therefore, the decomposition process of III<sup>3-</sup> is very similar to the CT induced anionic decompositions of other dioxetanones,<sup>30,31,59</sup> and its thermolysis and chemiexcitation can be also described by the GRCTIL mechanism. Meanwhile, similar results were found in the thermolysis of III<sup>-</sup>. (Figures S7 and S9) As Figure S6B shown,

the energies of non-CT  $S_0$  PEC (in the vicinities of reactant and product) is almost unaffected by DMSO, but the energies of CT region after TS is lower than that in gas phase.

No significant changes in key geometric structures and PECs of the III<sup>3-</sup> decomposition were found, when the reaction field is altered from gas phase to DMSO. Hence, considering the compromise between chemical accuracy and computational cost, we performed the TSH simulation in gas phase to confirm the GRCTIL mechanism and provide the dynamics information for  $III^{3-}$  decomposition. First, the average distributions of  $O_a O_{\rm b}$  and  $C_2-C_3$  at different time (Figure 8) clearly indicate a two-stage decay process. The fast  $O_a - O_b$  elongation takes place before 20 fs, while the  $C_2-C_3$  distance does not change obviously until the  $O_a - O_b$  distance reached ~2.5 Å. The further rupture of  $C_2 - C_3$  is observed after 40 fs. This indicates that the cleavages of  $O_a - O_b$  and  $C_2 - C_3$  are asynchronous, which is in line with the results in Figure 7C. Second, the evolution of average Mulliken charge population (Figure 9A) clearly displays the gradual CT from  $IV^{3-}$  to CO<sub>2</sub> (within ~20 fs), and a gradual BCT from CO<sub>2</sub> to  $IV^{3-}$  (20-60 fs), which respectively match with the asynchronous cleavages of  $O_a - O_b$ (0-20 fs) and  $C_2-C_3$  (>40 fs) (Figure 8) on the time scales and is also observed in Figure 7D. Finally, the transition stages of two electronic state were observed in the time evolution of average fractional occupations of  $S_0$  and  $S_1$  states (Figure 9B). In the first stage (stage 1), the average fractional occupation of S<sub>1</sub> increases from zero to approximately 0.7, implying the existence of a CI within 20 fs, where the obvious CT of  $IV^{3-} \rightarrow$  $CO_2$  and the  $O_a - O_b$  bond rupture takes place at the same time. Then the average fractional occupation of S<sub>1</sub> state decreases to 0.4 and becomes stable after 70 fs implying another CI in stage 2. The exactly matched time scales of the  $O_a - O_b/C_2 - C_3$  bond cleavage and the CT of  $IV^{3-} \rightarrow CO_2/BCT$  of  $CO_2 \rightarrow IV^{3-}$  is also consistent with above results in Figures 7C and 7D. Therefore, the two variation periods for  $S_1$  population could be attributed to two successive nonadiabatic transitions governed



**Figure 9.** Time evolution of the average Mulliken charge population of  $IV^{3-}$  and  $CO_2$  moieties (A), and the fractional occupations of the  $S_0$  and  $S_1$  states during the dynamics evolution of  $III^{3-}$  (B) (average over 100 trajectories) obtained by the TSH dynamics in gas phase at the SF-BHHLYP/6-31G(d,p) level.

by the two CIs. And two CIs in these two stages have been optimized by the SF-DFT method and correspond to the ruptures of  $O_a-C_b$  and  $C_2-C_3$  bonds, respectively (see Figure S8 for the geometric information on the two CIs). The same results were found in the TSH simulation of III<sup>-</sup> decomposition in gas phase (Figures S9–S11). Therefore, the results of TSH calculations and electronic structure computations support that the III<sup>3-</sup> decomposition follows the GRCTIL mechanism, and accomplishes nonadiabatic transition through the two CIs to emit light.

In addition, the quantum yield of  $IV^{3-}$  of the  $S_1$  state ( $\Phi s$ ) from the chemiexcitation in the III<sup>3-</sup> decomposition is also estimated according to the results of TSH simulation. For dioxetane-based CL or BL,  $\Phi$ s is the yield of electronic excitedstate products (e.g.,  $IV^{3-}$ ) formed during the decompositions of dioxetanones, i.e., the branch ratio of the S<sub>0</sub>- and S<sub>1</sub>-state after the nonadiabatic transition.<sup>75</sup>  $\Phi$ s and  $\Phi_{FL}$  (the fluorescence yield, i.e., the quantum yield of radiative decay of the excited product) determine the CL or BL yield ( $\Phi_{CL}$  or  $\Phi_{BL}$ ), that is  $\Phi_{\text{CL/BL}} = \Phi_{\text{S}} \times \Phi_{\text{FL}}$ . As shown in Figure 7B, at the end of TSH simulation (100 fs), III<sup>3-</sup> nearly decomposes completely, and 57 and 43 trajectories were found on the  $S_0$  and  $S_1$  states, respectively. Thus, the branch ratio of S<sub>1</sub> and S<sub>0</sub> is estimated to 43% in the chemiluminescent decomposition of III<sup>3-</sup> in gas phase, which is higher than that (34% in Figure S11B) in III<sup>-</sup> decomposition and is consistent in magnitude with the value (64%) estimated based on Landau–Zener formula (see details

in SI). Although there is no  $\Phi_{\rm FL}$  or  $\Phi_{\rm CL}$  of IV reported in experiments, the high  $\Phi_{\rm S}$  obtained by the TSH simulation suggests that the III<sup>3-</sup> decomposition could contribute to the high efficiency of *W. scintillans* BL. Combining the thermochemical data and dynamics simulation results, we suggest that III<sup>3-</sup> is more favorable and more efficient for *W. scintillans* BL.

#### 4. CONCLUSION

The luciferin of deep-sea squid, W. scintillans, has a universal core skeleton (ImPy) in about eight phyla of bioluminescent organisms. The BL mechanism proposed by Tsuji has not been confirmed by any strong theoretical or experimental evidence. The biggest challenge is to elucidate the mechanisms of two key steps, oxygenation of luciferin (I) and the thermolysis of the peroxide intermediate (III). The oxygenation process is a general step in almost all oxygen-dependent BL systems. The thermolysis of the parent skeleton of III, dioxetanone, is the key step to produce light emitter for most BL and organic CL. In this paper, the two steps were systematically investigated by reliable electronic structure computations at the DFT and TDDFT level, and the nonadiabatic molecular dynamics simulation with SF-DFT. Both the computed  $\Delta G_{\text{SET}}$  for the SET process and the estimated  $k_{\rm ET}$  based on Marcus theory demonstrate that the SET process is responsible for the oxygenation of  $I^{3-}$ . This SET process has been experimentally confirmed in the oxygenation of firefly luciferin. In DMSO, no matter whether  ${}^{1}[\mathbf{I}^{\bullet2-}\cdots\mathbf{O}_{2}^{\bullet-}]$  or  ${}^{3}[\mathbf{I}^{\bullet2-}\cdots\mathbf{O}_{2}^{\bullet-}]$  as the initial state, the C2-oxygentation is more energetically feasible than the C<sub>2</sub>-oxygentation.  ${}^1[I^{\bullet 2-} \cdots O_2^{\bullet -}]$  goes through four TSs to form III<sup>3-</sup> via nucleophilic addition followed by a biradical annihilation after TS2, while  ${}^{3}[I^{\bullet 2-}\cdots O_{2}^{\bullet -}]$  needs a ISC process via a S/T surface intersection to produce  $III^{3-}$  in C<sub>2</sub>oxygentation. The decompositions of III<sup>2-</sup> and III<sup>3-</sup> follow by entropic trapping and GRCTIL mechanisms, respectively, which is similar to the simple dioxetanones. The decomposition of III<sup>3-</sup> is energetically more feasible than III<sup>2-</sup>, III<sup>-</sup>, and III. The high quantum yield of  $IV^{3-}$  in the S<sub>1</sub> state ( $\Phi_S = 43\%$ ) indirectly supports that the III<sup>3-</sup>decomposition is the step to produce light emitter of the W. scintillans BL. The SEToxygenation and GRCTIL mechanism of W. scintillans BL in the current study could facilitate the understanding of most oxygen-dependent BL. The further study for the identification of light emitter of the W. scintillans BL is ongoing in our group.

#### ASSOCIATED CONTENT

#### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/jacs.6b09119.

The details for calculating the  $\Delta G_{\text{SET}}$  and  $k_{\text{ET}}$  for the SET process between I<sup>3-</sup> and <sup>3</sup>O<sub>2</sub>, and the estimation of  $\Phi_{\text{S}}$  in III<sup>3-</sup> decomposition based on Landau–Zener formula. The NOs used to represent the electronic states for I<sup>3-</sup> oxygenation and III<sup>2-</sup>/III<sup>3-</sup> decompositions; The S<sub>0</sub> PECs of the C<sub>2</sub>- and C<sub>3</sub>-oxygenation of I<sup>-</sup> ( $-\text{OSO}_3\text{H}$ ) in DMSO computed at the UCAM-B3LYP/6-31G(d,p) level. The S<sub>0</sub> PECs of the C<sub>2</sub>- and C<sub>3</sub>-oxygenation of I<sup>3-</sup> in DMSO computed at the UCAM-B3LYP/6-31G(d,p) level with dispersion. The comparison of S<sub>0</sub> PECs of the III<sup>2-</sup>/III<sup>3-</sup> ( $-\text{OSO}_3^-$ ) decomposition in gas phase and DMSO calculated at the UCAM-B3LYP/6-31G(d,p) level; The S<sub>0</sub> and approximated diabatic PECs, and the variation of the key geometric parameters and Mulliken charge population in III/III<sup>-</sup>( $-OSO_3H$ ) decompositions calculated by (TD) UCAM-B3LYP/6-31G(d,p); The time evolution of detailed two-dimensional distribution of the C<sub>2</sub>-C<sub>3</sub> and O<sub>a</sub>-O<sub>b</sub> bond distances, the average Mulliken charge populations of IV/IV<sup>-</sup>( $-OSO_3H$ ) and CO<sub>2</sub> moieties, and the fractional occupation of S<sub>0</sub> and S<sub>1</sub> states obtained by TSH simulation at the SF-BHHLYP/ 6-31G(d,p) level; Cartesian coordinates for the RC, Int, and TS optimized at the (TD) CAM-B3LYP/6-31G(d,p) level, and CIs optimized at the SF-BHHLYP/ 6-31G(d,p) level; (PDF)

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#### Notes

The authors declare no competing financial interest.

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#### REFERENCES

(1) Bianchi, A.; Dufort, S.; Lux, F.; Fortin, P.-Y.; Tassali, N.; Tillement, O.; Coll, J.-L.; Crémillieux, Y. Proc. Natl. Acad. Sci. U. S. A. **2014**, 111, 9247–9252.

(2) Schaub, F. X.; Reza, M. S.; Flaveny, C. A.; Li, W.; Musicant, A. M.; Hoxha, S.; Guo, M.; Cleveland, J. L.; Amelio, A. L. *Cancer Res.* **2015**, 75, 5023–5033.

(3) Wang, F.; Wang, Z.; Hida, N.; Kiesewetter, D. O.; Ma, Y.; Yang, K.; Rong, P.; Liang, J.; Tian, J.; Niu, G.; Chen, X. *Proc. Natl. Acad. Sci.* U. S. A. **2014**, *111*, 5165–5170.

(4) Markova, S. V.; Vysotski, E. S. Biochemistry (Moscow) 2015, 80, 714–732.

(5) Purtov, K. V.; Petushkov, V. N.; Baranov, M. S.; Mineev, K. S.; Rodionova, N. S.; Kaskova, Z. M.; Tsarkova, A. S.; Petunin, A. I.; Bondar, V. S.; Rodicheva, E. K.; Medvedeva, S. E.; Oba, Y.; Oba, Y.; Arseniev, A. S.; Lukyanov, S.; Gitelson, J. I.; Yampolsky, I. V. Angew. Chem., Int. Ed. **2015**, 54, 8124–8128.

(6) Dubinnyi, M. A.; Kaskova, Z. M.; Rodionova, N. S.; Baranov, M. S.; Gorokhovatsky, A. Y.; Kotlobay, A.; Solntsev, K. M.; Tsarkova, A. S.; Petushkov, V. N.; Yampolsky, I. V. *Angew. Chem., Int. Ed.* **2015**, *54*, 7065–7067.

(7) Haddock, S. H. D.; Moline, M. A.; Case, J. F. Annu. Rev. Mar. Sci. 2010, 2, 443–493.

- (8) Widder, E. A. Science 2010, 328, 704-708.
- (9) Tsuji, F. I. Proc. Natl. Acad. Sci. U. S. A. 1985, 82, 4629-4632.

(10) Tsuji, F. I. Biochem. Biophys. Res. Commun. 2005, 338, 250–253.
(11) Tsuji, F. I. Biochim. Biophys. Acta, Biomembr. 2002, 1564, 189–

197. (12) Hayashi, K.; Kawai, Y. L.; Yura, K.; Yoshida, M.-a.; Ogura, A.; Hata, K.; Nakabayashi, K.; Okamura, K. *Mitochondrial DNA* **2016**, 1–

(13) DeLoney-Marino, C. R.; Wolfe, A. J.; Visick, K. L. Appl. Environ. Microbiol. 2003, 69, 7527–7530.

(14) Yazzie, N.; Salazar, K. A.; Castillo, M. G. Fish Shellfish Immunol. 2015, 44, 342–355.

(15) Tong, D.; Rozas, N. S.; Oakley, T. H.; Mitchell, J.; Colley, N. J.; McFall-Ngai, M. J. *Proc. Natl. Acad. Sci. U. S. A.* **2009**, *106*, 9836– 9841.

- (16) Jones, B. W.; Nishiguchi, M. K. Mar. Biol. 2004, 144, 1151–1155.
- (17) Dilly, P. N.; Herring, P. J. J. Zool. 1978, 186, 47-59.
- (18) Tsuji, F. I.; Leisman, G. B. Proc. Natl. Acad. Sci. U. S. A. 1981, 78, 6719-6723.
- (19) Fujii, T.; Ahn, J.-Y.; Kuse, M.; Mori, H.; Matsuda, T.; Isobe, M. *Biochem. Biophys. Res. Commun.* **2002**, *293*, 874–879.
- (20) Chou, C.-M.; Tung, Y.-W.; Isobe, M. Bioorg. Med. Chem. 2014, 22, 4177–4188.
- (21) Isobe, M.; Kuse, M.; Tani, N.; Fujii, T.; Matsuda, T. Proc. Jpn. Acad., Ser. B **2008**, 84, 386–392.
- (22) Kuse, M.; Isobe, M. Tetrahedron 2000, 56, 2629-2639.
- (23) Kuse, M. Biosci., Biotechnol., Biochem. 2014, 78, 731-736.
- (24) Goto, T.; Iio, H.; Inoue, S.; Kakoi, H. Tetrahedron Lett. 1974, 15, 2321–2324.

(25) Inoue, S.; Sugiura, S.; Kakoi, H.; Hasizume, K.; Goto, T.; Iio, H. *Chem. Lett.* **1975**, *4*, 141–144.

- (26) Inoue, S.; Kakoi, H.; Goto, T. Tetrahedron Lett. **1976**, 17, 2971–2974.
- (27) Haddock, S. H. D.; Rivers, T. J.; Robison, B. H. Proc. Natl. Acad. Sci. U. S. A. 2001, 98, 11148–11151.
- (28) Jiang, T.; Du, L.; Li, M. Photochem. Photobiol. Sci. 2016, 15, 466–480.
- (29) Teranishi, K.; Shimomura, O. Biochim. Biophys. Acta, Gen. Subj. 2008, 1780, 784–792.
- (30) Yue, L.; Liu, Y.-J.; Fang, W.-H. J. Am. Chem. Soc. 2012, 134, 11632-11639.
- (31) Ding, B.-W.; Naumov, P.; Liu, Y.-J. J. Chem. Theory Comput. 2015, 11, 591–599.
- (32) Chen, S.-F.; Ferré, N.; Liu, Y.-J. Chem. Eur. J. 2013, 19, 8466–8472.
- (33) Liu, Z.-J.; Stepanyuk, G. A.; Vysotski, E. S.; Lee, J.; Markova, S. V.; Malikova, N. P.; Wang, B.-C. *Proc. Natl. Acad. Sci. U. S. A.* **2006**, 103, 2570–2575.
- (34) Hou, C.; Liu, Y.-J.; Ferré, N.; Fang, W.-H. Chem. Eur. J. 2014, 20, 7979-7986.
- (35) McCapra, F.; Chang, Y. C. Chem. Commun. (London) 1967, 1011–1012.
- (36) McCapra, F. Acc. Chem. Res. 1976, 9, 201-208.
- (37) Shimomura, O.; Johnson, F. H. Biochem. Biophys. Res. Commun. 1971, 44, 340–346.
- (38) Kazuo Hori, M. J. C. Proc. Natl. Acad. Sci. U.S.A. 1973, 70, 4.
- (39) Ohmiya, Y.; Hirano, T. Chem. Biol. 1996, 3, 337-347.
- (40) Adamczyk, M.; Akireddy, S. R.; Johnson, D. D.; Mattingly, P. G.; Pan, Y.; Reddy, R. E. *Tetrahedron* **2003**, *59*, 8129–8142.
- (41) Teranishi, K.; Ueda, K.; Nakao, H.; Hisamatsu, M.; Yamada, T. Tetrahedron Lett. **1994**, 35, 8181–8184.
- (42) Goto, T. Pure Appl. Chem. 1968, 17, 421-442.
- (43) Nakamura, H.; Wu, C.; Inouye, S.; Murai, A. J. Am. Chem. Soc. 2001, 123, 1523–1524.
- (44) Isobe, M.; Takahashi, H.; Usami, K.; Hattori, M.; Nishigohri, Y. Pure Appl. Chem. **1994**, *66*, 765.
- (45) Ken Fujimori, S. H. N.; Katsuhisa, A.; Motohiro, M. J. Chem. Soc. Perkin Trans. 2 1993, 2, 5.
- (46) Isobe, H.; Yamanaka, S.; Kuramitsu, S.; Yamaguchi, K. J. Am. Chem. Soc. 2008, 130, 132–149.
- (47) Usami, K.; Isobe, M. Tetrahedron 1996, 52, 12061-12090.
- (48) Head, J. F.; Inouye, S.; Teranishi, K.; Shimomura, O. Nature 2000, 405, 372–376.
- (49) Kondo, H.; Igarashi, T.; Maki, S.; Niwa, H.; Ikeda, H.; Hirano, T. *Tetrahedron Lett.* **2005**, *46*, 7701–7704.
- (50) Liu, Z.-J.; Stepanyuk, G. A.; Vysotski, E. S.; Lee, J.; Markova, S. V.; Malikova, N. P.; Wang, B.-C. *Proc. Natl. Acad. Sci. U. S. A.* **2006**, 103, 2570–2575.
- (51) Visitsatthawong, S.; Chenprakhon, P.; Chaiyen, P.; Surawatanawong, P. J. Am. Chem. Soc. 2015, 137, 9363-9374.

<sup>(52)</sup> Branchini, B. R.; Behney, C. E.; Southworth, T. L.; Fontaine, D. M.; Gulick, A. M.; Vinyard, D. J.; Brudvig, G. W. J. Am. Chem. Soc. 2015, 137, 7592–7595.

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(53) Eremeeva, E. V.; Natashin, P. V.; Song, L.; Zhou, Y.; van Berkel,

W. J. H.; Liu, Z.-J.; Vysotski, E. S. ChemBioChem 2013, 14, 739–745.
 (54) Chen, S.-F.; Navizet, I.; Roca-Sanjuán, D.; Lindh, R.; Liu, Y.-J.;

- Ferré, N. J. Chem. Theory Comput. 2012, 8, 2796–2807. (55) Chen, S.; Navizet, I.; Lindh, R.; Liu, Y.; Ferré, N. J. Phys. Chem.
- B **2014**, *118*, 2896–2903.

(56) Natashin, P. V.; Ding, W.; Eremeeva, E. V.; Markova, S. V.; Lee, J.; Vysotski, E. S.; Liu, Z.-J. Acta Crystallogr., Sect. D: Biol. Crystallogr. 2014, 70, 720–732.

(57) Yoshihiro, O.; Satoshi, K.; Mitsuhiro, N.; Haruki, N. Bull. Chem. Soc. Jpn. 2005, 78, 1197–1205.

(58) White, E. H.; Zafiriou, O.; Kagi, H. H.; Hill, J. H. M. J. Am. Chem. Soc. 1964, 86, 940-941.

(59) Yue, L.; Liu, Y.-J. J. Chem. Theory Comput. 2013, 9, 2300-2312.

(60) King, D. W.; Cooper, W. J.; Rusak, S. A.; Peake, B. M.; Kiddle, J. J.; O'Sullivan, D. W.; Melamed, M. L.; Morgan, C. R.; Theberge, S. M. *Anal. Chem.* **2007**, *79*, 4169–4176.

(61) Vico, L. D.; Liu, Y.-J.; Krogh, J. W.; Lindh, R. J. Phys. Chem. A 2007, 111, 8013–8019.

(62) Yue, L.; Roca-Sanjuán, D.; Lindh, R.; Ferré, N.; Liu, Y.-J. J. Chem. Theory Comput. 2012, 8, 4359-4363.

(63) Liu, F.; Liu, Y.-J.; Vico, L. D.; Lindh, R. J. Am. Chem. Soc. 2009, 131, 6181–6188.

(64) Liu, F.; Liu, Y.-J.; Vico, L. D.; Lindh, R. Chem. Phys. Lett. 2009, 484, 69–75.

(65) Yanai, T.; Tew, D. P.; Handy, N. C. Chem. Phys. Lett. 2004, 393, 51–57.

(66) Marcus, R. A. Rev. Mod. Phys. 1993, 65, 599.

- (67) Krylov, A. I. Chem. Phys. Lett. 2001, 338, 375-384.
- (68) Krylov, A. I.; Sherrill, C. D. J. Chem. Phys. 2002, 116, 3194–3203.

(69) Casanova, D.; Head-Gordon, M. J. Chem. Phys. 2008, 129, 064104.

(70) Hirata, S.; Head-Gordon, M. Chem. Phys. Lett. **1999**, 314, 291–299.

(71) Huix-Rotllant, M.; Natarajan, B.; Ipatov, A.; Muhavini Wawire, C.; Deutsch, T.; Casida, M. E. Phys. Chem. Chem. Phys. 2010, 12, 12811–12825.

(72) Liu, F.; Gan, Z.; Shao, Y.; Hsu, C.-P.; Dreuw, A.; Head-Gordon, M.; Miller, B. T.; Brooks, B. R.; Yu, J.-G.; Furlani, T. R.; Kong, J. *Mol. Phys.* **2010**, *108*, 2791–2800.

(73) Harabuchi, Y.; Maeda, S.; Taketsugu, T.; Minezawa, N.; Morokuma, K. J. Chem. Theory Comput. **2013**, *9*, 4116–4123.

(74) Minezawa, N.; Gordon, M. S. J. Phys. Chem. A 2009, 113, 12749-12753.

(75) Yue, L.; Lan, Z.; Liu, Y.-J. J. Phys. Chem. Lett. 2015, 6, 540–548.
(76) Pittner, J.; Lischka, H.; Barbatti, M. Chem. Phys. 2009, 356,

147–152. (77) Werner, U.; Mitrić, R.; Suzuki, T.; Bonačić-Koutecký, V. *Chem. Phys.* **2008**, 349, 319–324.

- (78) Verlet, L. Phys. Rev. **1967**, 159, 98–103.
- (79) Verlet, L. Phys. Rev. **1968**, 165, 201–214.

(80) Swope, W. C.; Andersen, H. C.; Berens, P. H.; Wilson, K. R. J. Chem. Phys. **1982**, 76, 637-649.

(81) Sellner, B.; Barbatti, M.; Lischka, H. J. Chem. Phys. 2009, 131, 024312.

(82) Liu, Y.-J.; Tian, Y.-C.; Fang, W.-H. J. Chem. Phys. 2010, 132, 014306.

(83) Landau, L. D. Phys. Z. Sowjetunion 1932, 2, 46-51.

(84) Zener, C. Proc. R. Soc. London, Ser. A 1932, 137, 696-702.

(85) Hehre, W. J.; Ditchfield, R.; Pople, J. A. J. Chem. Phys. 1972, 56, 2257-2261.

(86) Hariharan, P. C.; Pople, J. A. Theoretical Chemistry Accounts: Theory, Computation, and Modeling (Theoretica Chimica Acta) 1973, 28, 213–222.

(87) Frisch, M. J.; Trucks, G. W.; Schlegel, H. B.; Scuseria, G. E.; Robb, M. A.; Cheeseman, J. R.; Scalmani, G.; Barone, V.; Mennucci, B.; Petersson, G. A.; Nakatsuji, H.; Caricato, M.; Li, X.; Hratchian, H. P.; Izmaylov, A. F.; Bloino, J.; Zheng, G.; Sonnenberg, J. L.; Hada, M.; Ehara, M.; Toyota, K.; Fukuda, R.; Hasegawa, J.; Ishida, M.; Nakajima, T.; Honda, Y.; Kitao, O.; Nakai, H.; Vreven, T.; Montgomery, J. A., Jr.; Peralta, J. E.; Ogliaro, F.; Bearpark, M.; Heyd, J. J.; Brothers, E.; Kudin, K. N.; Staroverov, V. N.; Kobayashi, R.; Normand, J.; Raghavachari, K.; Rendell, A.; Burant, J. C.; Iyengar, S. S.; Tomasi, J.; Cossi, M.; Rega, N.; Millam, J. M.; Klene, M.; Knox, J. E.; Cross, J. B.; Bakken, V.; Adamo, C.; Jaramillo, J.; Gomperts, R.; Stratmann, R. E.; Yazyev, O.; Austin, A. J.; Cammi, R.; Pomelli, C.; Ochterski, J. W.; Martin, R. L.; Morokuma, K.; Zakrzewski, V. G.; Voth, G. A.; Salvador, P.; Dannenberg, J. J.; Dapprich, S.; Daniels, A. D.; Farkas, O.; Foresman, J. B.; Ortiz, J. V.; Cioslowski, J.; Fox, D. J. *Gaussian 09*, revision A.02; Gaussian, Inc.: Wallingford, CT, 2009.

(88) Barbatti, M.; Granucci, G.; Persico, M.; Ruckenbauer, M.; Vazdar, M.; Eckert-Maksić, M.; Lischka, H. J. Photochem. Photobiol., A 2007, 190, 228–240.

(89) Barbatti, M.; Granucci, G.; Ruckenbauer, M.; Pittner, J.; Persico, M.; Lischka, H. NEWTON-X: A Package for Newtonian Dynamics Close to the Crossing Seam, version 1.3; 2012; http://www.newtonx.org.

(90) Schmidt, M. W.; Baldridge, K. K.; Boatz, J. A.; Elbert, S. T.; Gordon, M. S.; Jensen, J. H.; Koseki, S.; Matsunaga, N.; Nguyen, K. A.; Su, S.; Windus, T. L.; Dupuis, M.; Montgomery, J. A. J. Comput. Chem. **1993**, *14*, 1347–1363.

- (91) Noyes, R. M. J. Am. Chem. Soc. 1955, 77, 2042-2045.
- (92) Noyes, R. M. J. Am. Chem. Soc. 1956, 78, 5486-5490.
- (93) Winstein, S.; Clippinger, E.; Fainberg, A. H.; Robinson, G. C. J. Am. Chem. Soc. 1954, 76, 2597–2598.
- (94) Kavarnos, G. J.; Turro, N. J. Chem. Rev. 1986, 86, 401-449.
- (95) Werner, H. J.; Staerk, H.; Weller, A. J. Chem. Phys. 1978, 68, 2419-2426.
- (96) Steiner, U. E.; Ulrich, T. Chem. Rev. 1989, 89, 51-147.
- (97) El-Sayed, M. A. J. Chem. Phys. 1963, 38, 2834-2838.